

### INSTRUCTIONS

- Show ALL your work. Insufficient justification will result in a loss of points.
- Electronic calculators, computers, etc. are not permitted.

1. (20 pts.) Use the standard definition of the angular momentum operator,

$$\hat{\mathbf{L}} = \hat{\mathbf{r}} \times \hat{\mathbf{p}}, \quad (1)$$

where  $\hat{\mathbf{r}} = (\hat{x}, \hat{y}, \hat{z})$  and  $\hat{\mathbf{p}} = (\hat{p}_x, \hat{p}_y, \hat{p}_z)$  are the coordinate and momentum operators,<sup>#</sup> and the commutation relationships,

$$[\hat{x}_k, \hat{x}_\ell] = 0, \quad [\hat{p}_k, \hat{p}_\ell] = 0, \quad [\hat{x}_k, \hat{p}_\ell] = i\hbar\delta_{k,\ell} \quad (k, \ell = 1, 2, 3), \quad (2)$$

where  $\hat{x}_1 \equiv \hat{x}$ ,  $\hat{x}_2 \equiv \hat{y}$ ,  $\hat{x}_3 \equiv \hat{z}$ ,  $\hat{p}_1 \equiv \hat{p}_x$ ,  $\hat{p}_2 \equiv \hat{p}_y$ , and  $\hat{p}_3 \equiv \hat{p}_z$ , to prove that

$$[\hat{L}_x, \hat{L}_y] = i\hbar\hat{L}_z. \quad (3)$$

The derivation must be purely algebraic (i.e., without any reference to a particular form of the quantum-mechanical operators in the coordinate representation).

2. (15pts.) Let  $\hat{\mathbf{J}} = (\hat{J}_x, \hat{J}_y, \hat{J}_z)$  be the angular momentum operator. Use the commutation relationships,

$$[\hat{J}_x, \hat{J}_y] = i\hbar\hat{J}_z, \quad (4)$$

$$[\hat{J}_y, \hat{J}_z] = i\hbar\hat{J}_x, \quad (5)$$

$$[\hat{J}_z, \hat{J}_x] = i\hbar\hat{J}_y, \quad (6)$$

to prove that

$$[\hat{J}^2, \hat{J}_z] = 0, \quad (7)$$

$$[\hat{J}_+, \hat{J}_-] = 2\hbar\hat{J}_z, \quad (8)$$

where  $\hat{J}^2 = \hat{J}_x^2 + \hat{J}_y^2 + \hat{J}_z^2$  and  $\hat{J}_\pm = \hat{J}_x \pm i\hat{J}_y$ .

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<sup>#)</sup> Bold symbols designate vector (or tensor) quantities.

3. (15 pts.) Let  $|jm\rangle$  be an eigenstate of  $\hat{J}^2 = \hat{J}_x^2 + \hat{J}_y^2 + \hat{J}_z^2$  and  $\hat{J}_z$ , where  $\hat{J}_x$ ,  $\hat{J}_y$ , and  $\hat{J}_z$  are the Cartesian components of the angular momentum operator  $\hat{\mathbf{J}}$ . Use the above Eqs. (4)–(6) to show that the expectation values of  $\hat{J}_x$  and  $\hat{J}_y$  in the state  $|jm\rangle$  are zero, i.e.,

$$\langle jm|\hat{J}_x|jm\rangle = \langle jm|\hat{J}_y|jm\rangle = 0. \quad (9)$$

*Hint:* Use Eqs. (4)–(6) to express  $\hat{J}_x$  and  $\hat{J}_y$  in terms of the commutators of the components of the angular momentum operator.

4. (25 pts.) Let  $\alpha(i)$  and  $\beta(i)$ ,  $i = 1, 2$ , be the spin-up and spin-down spin functions of electrons 1 and 2, respectively. Let  $\hat{\mathbf{S}}_1 = (\hat{S}_{1,x}, \hat{S}_{1,y}, \hat{S}_{1,z})$  and  $\hat{\mathbf{S}}_2 = (\hat{S}_{2,x}, \hat{S}_{2,y}, \hat{S}_{2,z})$  be the spin angular momentum operators corresponding to electrons 1 and 2, respectively, and let  $\hat{\mathbf{S}} = \hat{\mathbf{S}}_1 + \hat{\mathbf{S}}_2$  be the operator representing the total spin of a two-electron system consisting of electrons 1 and 2. Use the values of Clebsch-Gordan coefficients  $\langle j_1 m_1, j_2 m_2 | jm \rangle$  for  $j_1 = j_2 = \frac{1}{2}$  provided at the end of this document to represent the eigenstates  $|sm\rangle$  of  $\hat{S}^2 = (\hat{\mathbf{S}}_1 + \hat{\mathbf{S}}_2)^2$  and  $\hat{S}_z = \hat{S}_{1,z} + \hat{S}_{2,z}$  as linear combinations of the two-electron spin functions  $\alpha(1)\alpha(2)$ ,  $\alpha(1)\beta(2)$ ,  $\beta(1)\alpha(2)$ ,  $\beta(1)\beta(2)$ . Report your results in the form of the following table:

eigenstate $ sm\rangle$ of $\hat{S}^2$ and $\hat{S}_z$	eigenvalue of $\hat{S}^2$	eigenvalue of $\hat{S}_z$
...	...	...

Show that the  $|sm\rangle$  eigenstates of  $\hat{S}^2$  and  $\hat{S}_z$  resulting from your calculations form an orthonormal set.

5. (25 pts.) (a) The  $\hat{J}_\pm = \hat{J}_x \pm i\hat{J}_y$  operators are called the angular momentum raising (+) and lowering (–) operators (as usual,  $\hat{J}_x$ ,  $\hat{J}_y$ , and  $\hat{J}_z$  are the Cartesian components of the angular momentum operator  $\hat{\mathbf{J}}$ ). Show that

$$\hat{J}^2 = \frac{1}{2}(\hat{J}_+\hat{J}_- + \hat{J}_-\hat{J}_+) + \hat{J}_z^2. \quad (10)$$

(b) Consider the three-electron system consisting of electrons 1, 2, and 3. Let  $\alpha(i)$  and  $\beta(i)$ ,  $i = 1, 2, 3$ , be the spin-up and spin-down spin functions of electrons 1, 2, and 3, respectively. Let  $\hat{\mathbf{S}}_i = (\hat{S}_{i,x}, \hat{S}_{i,y}, \hat{S}_{i,z})$ ,  $i = 1, 2, 3$ , be the spin angular momentum operators corresponding to electrons 1, 2, and 3, respectively. It can be shown that

$$\hat{S}_{i,+}\alpha(i) = 0, \quad \hat{S}_{i,-}\alpha(i) = \hbar\beta(i), \quad \hat{S}_{i,+}\beta(i) = \hbar\alpha(i), \quad \hat{S}_{i,-}\beta(i) = 0, \quad (11)$$

where  $\hat{S}_{i,\pm} = \hat{S}_{i,x} \pm i\hat{S}_{i,y}$  are the spin angular momentum raising and lowering operators ( $i = 1, 2, 3$ ). Use Eqs. (10) and (11) to show that the following function:

$$u(1, 2, 3) = \frac{1}{\sqrt{6}}[2\beta(1)\alpha(2)\alpha(3) - \alpha(1)\alpha(2)\beta(3) - \alpha(1)\beta(2)\alpha(3)] \equiv \frac{1}{\sqrt{6}}[2\beta\alpha\alpha - \alpha\alpha\beta - \alpha\beta\alpha], \quad (12)$$

is an eigenstate of the total spin operator  $\hat{S}^2 = (\hat{\mathbf{S}}_1 + \hat{\mathbf{S}}_2 + \hat{\mathbf{S}}_3)^2$ . What is the corresponding eigenvalue?

The Clebsch-Gordan coefficients  $\langle j_1 m_1, j_2 m_2 | j m \rangle$  for  $j_1 = j_2 = \frac{1}{2}$ .

		$j = 1$ $m = 1$	$j = 1$ $m = 0$	$j = 0$ $m = 0$	$j = 1$ $m = -1$
$m_1 = \frac{1}{2}$	$m_2 = \frac{1}{2}$	1	0	0	0
$m_1 = \frac{1}{2}$	$m_2 = -\frac{1}{2}$	0	$\frac{1}{\sqrt{2}}$	$\frac{1}{\sqrt{2}}$	0
$m_1 = -\frac{1}{2}$	$m_2 = \frac{1}{2}$	0	$\frac{1}{\sqrt{2}}$	$-\frac{1}{\sqrt{2}}$	0
$m_1 = -\frac{1}{2}$	$m_2 = -\frac{1}{2}$	0	0	0	1

### Other potentially useful formulas

- $[\hat{F}_1 \hat{F}_2, \hat{G}] = \hat{F}_1 [\hat{F}_2, \hat{G}] + [\hat{F}_1, \hat{G}] \hat{F}_2$  ( $[\ , \ ]$  is the commutator)
- $[F_1 F_2, G_1 G_2] = F_1 G_1 [F_2, G_2] + F_1 [F_2, G_1] G_2 + G_1 [F_1, G_2] F_2 + [F_1, G_1] G_2 F_2$ .